

Locked Constraint Satisfaction Problems

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We introduce and study the random “locked” constraint satisfaction problems. When increasing the density of constraints, they display a broad “clustered” phase in which the space of solutions is divided into many *isolated* points. While the phase diagram can be found easily, these problems, in their clustered phase, are extremely hard from the algorithmic point of view: the best known algorithms all fail to find solutions. We thus propose new benchmarks of really hard optimization problems and provide insight into the origin of their typical hardness.

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Constraint satisfaction problems (CSPs) are one of the main building blocks of complex systems studied in computer science, information theory, and statistical physics. Their wide range of applicability arises from their very general nature: given a set of N discrete variables subject to M constraints, the CSP consists in deciding whether there exists an assignment of variables which satisfies simultaneously all the constraints. In computer science, CSPs are at the core of computational complexity studies: the satisfiability of boolean formulas is the canonical example of an intrinsically hard, NP-complete, problem [1]. In information theory, error-correcting codes also rely on CSPs. The transmitted information is encoded into a code word satisfying a set of constraints so that information may be retrieved after transmission through a noisy channel, using the knowledge of the constraints. Many other practical problems in scheduling a collection of tasks or in hardware and software verification and testing are viewed as CSPs. In statistical physics, the interest in CSPs stems from their close relation with theory of spin glasses. Answering if frustration is avoidable in a system is a first, and sometimes highly nontrivial, step in understanding its low-temperature behavior.

Methods of statistical physics provide powerful tools to study statistical properties of CSPs [2,3]. The mean field approach is known to be exact if the underlying graph of constraints [4] is either fully connected or locally treelike. It also has algorithmic, and practical, consequences: in contrast with the usual situation in physics, CSPs on a locally treelike graph are used in practice, for instance in low density parity check codes [5], which are among the best error-correcting codes around.

Many CSPs are NP-complete. Nevertheless, large classes of instances can be easy to solve. It is one of the main goals of theoretical computer science to understand why some instances are harder than others, where the hardness comes from and how to avoid it, beat it, or use it. The random K -satisfiability (K -SAT) problem where clauses are chosen uniformly at random between all possible ones has played a prominent role in approaching this

goal. In random K -SAT, there exists a sharp satisfiability threshold. This is a phase transition point separating a “SAT” phase with low density of constraints where instances are almost always satisfiable, from an “UNSAT” phase where, with high probability, there is no solution to the CSP [6,7]. The hardest instances lie near to this threshold [8,9]. The main insight came from statistical physics studies [10–16] which allow to describe the structure of the space of solution of the random K -SAT problem. The most interesting result is the existence of an intermediate “clustered” phase, just below the SAT-UNSAT threshold, where the space of solutions splits into well-separated clusters. A major open question consists in understanding if and how the existence of clusters makes the problem harder. The survey propagation algorithm, which explicitly takes into account the clusters, is the best known solver very close to the SAT-UNSAT threshold [13], but some local search algorithms also perform well inside the clustered phase [17,18]. Another proposition, put forward in [15], is that solutions in clusters with frozen variables, taking the same value in the whole cluster, are hard to find. It was shown in [19] that, even if solutions belonging to clusters without frozen variables are exponentially rare, some message passing algorithms may be able to find them.

In this Letter, we introduce and study a broad class of CSPs which are extremely frozen problems: all the clusters consist of a single configuration; thus, all the variables are frozen in every cluster. We show that these problems are extremely difficult from an algorithmic point of view: all the best known algorithms fail to solve them in this clustered phase. At the same time, the description of their phase diagram can be carried out in details with relatively simple statistical physics methods.

Definition.—We define an *occupation CSP* over N binary variables, $s_1, \dots, s_N \in \{0, 1\}$ as follows: each constraint a connects to K randomly chosen variables, and its status depends on the sum r of these variables. The constraint is characterized by a $(K + 1)$ component vector $A = (A_0 A_1 \dots A_K)$, with $A_r \in \{0, 1\}$: it is satisfied if and only if $A_r = 1$. We shall study here homogeneous models

in which all constraints connect to the same number K of variables, and are characterized by the same vector A . According to [20], the occupation CSPs are NP-complete if $K > 2$, $A_0 = A_K = 0$, and A is not a parity check. The *locked occupation problems* (LOP) are occupation CSPs satisfying two conditions: (a) $\forall i = 0, \dots, K-1$ the product $A_i A_{i+1} = 0$, (b) all variables are present in at least two constraints. Simple examples of LOPs are positive 1-in-3 satisfiability [21], $A = 0100$, or parity checks [5], $A = 01010$, on graphs without leaves. In order to go from one solution (satisfying assignment) of a LOP to another one, it is necessary to flip at least a closed loop of variables in the factor graph representation of [4]. This stays at the root of the crucial property that clusters are pointlike and separated by an extensive distance when the density of constraints is large enough (above l_d).

In order to fully characterize a random LOP ensemble, one needs to define the degree distribution of variables. We will study here two ensembles. The regular ensemble, where every variable appears in exactly L constraints, and the truncated Poisson ensemble with degree distribution $Q(0) = Q(1) = 0$, $Q(l) = e^{-c} c^l / l! [1 - (1+c)e^{-c}]$, $l \geq 2$, and average connectivity $\bar{l} = c(1 - e^{-c}) / [1 - (1+c)e^{-c}]$.

Phase diagram.—Denoting by a, b, \dots , the indices of constraints and i, j, \dots those of variables, the belief propagation (BP) equations [22] are given by

$$\psi_{s_i}^{a \rightarrow i} = \frac{1}{Z^{a \rightarrow i}} \sum_{\{s_j\}} \delta(A_{s_i + \sum_j s_j} - 1) \prod_{j \in \partial a - i} \chi_{s_j}^{j \rightarrow a}, \quad (1)$$

$$\chi_{s_j}^{j \rightarrow a} = \frac{1}{Z^{j \rightarrow a}} \prod_{b \in \partial j - a} \psi_{s_j}^{b \rightarrow j}, \quad (2)$$

where ∂a are all the variables appearing in constraint a , and ∂i all the constraints in which variable i appears. $\chi_{s_j}^{j \rightarrow a}$ is the probability that spin j takes value s_j when a was removed from the graph, and Z are normalization constants. The BP entropy (the logarithm of number of configuration satisfying all constraints, divided by N) is

$$s = \frac{1}{N} \sum_a \log(Z^{a+\partial a}) - \frac{1}{N} \sum_i (l_i - 1) \log(Z^i), \quad (3)$$

where

$$Z^{a+\partial a} = \sum_{\{s_i\}} \delta(A_{\sum_i s_i} - 1) \prod_{i \in \partial a} \left(\prod_{b \in \partial i - a} \psi_{s_i}^{b \rightarrow i} \right), \quad (4)$$

$$Z^i = \prod_{a \in \partial i} \psi_0^{a \rightarrow i} + \prod_{a \in \partial i} \psi_1^{a \rightarrow i}. \quad (5)$$

In order to find a fixed point of Eqs. (1) and (2) and compute the quenched average of the entropy, we use the population dynamics technique [2], with population sizes of order 10^4 to 10^5 . It turns out that this procedure always converges to the same fixed point.

The phase diagram of LOPs is much simpler to analyze than the one of general CSPs, and can be deduced purely from the BP analysis. This is due to the fact that, in the clustered phase, every cluster reduces to a single isolated configuration. The survey propagation (SP) equations [13] are then greatly simplified. Their iteration either leads to a trivial fixed point, where every variable is in the so called “joker” state [23], or to a fixed point where no variable is in the joker state. In this second case, the SP equations reduce to the BP Eqs. (1) and (2), and the complexity function (logarithm of number of clusters) is equal to the entropy (3), in agreement with the pointlike nature of clusters. The clustered phase is then identified from the iterative stability of this second, nontrivial fixed point. It is iteratively stable when the average connectivity is above a threshold: $\bar{l} > l_d$, while the regime $\bar{l} < l_d$ corresponds to a “liquid” phase. The intuitive difference between the two phases is that in the clustered phase, one has to flip an extensive number of variables to go from one solution to another, while in the liquid phase, the addition of any infinitesimal temperature is enough to be able to connect all solutions.

The satisfiability threshold l_s is defined as follows: If the average connectivity is $\bar{l} < l_s$, then a satisfying assignment almost surely exists (in $N \rightarrow \infty$), and if $\bar{l} > l_s$, then there is almost surely no satisfying configuration. In LOPs, we can find l_s as the average connectivity at which the RS entropy (3) becomes zero. Table I gives the values of clustering and satisfiability thresholds for the nontrivial LOPs with $K \leq 5$.

When a LOP is symmetric, i.e., $A_r = A_{K-r}$ for all $r = 0, \dots, K$, and this 0–1 symmetry is not spontaneously broken, the satisfiability threshold can be computed rigorously using the 1st and the 2nd moment methods: The annealed entropy $\langle Z \rangle = \exp(Ns_{\text{ann}})$ is

$$s_{\text{ann}}(\bar{l}) = \log 2 + \frac{\bar{l}}{K} \log \left[2^{-K} \sum_{r=0}^K \delta(A_r - 1) \binom{K}{r} \right]. \quad (6)$$

TABLE I. The clustering l_d and satisfiability l_s thresholds in the locked occupation problems for $K \leq 5$ in the truncated Poisson ensemble. In the regular ensemble, L_s is the first unsatisfiable or critical connectivity; the first clustered case is $L_d = 3$. The error bars originate in the statistical nature of the population dynamics technique. Symmetric LOPs where the satisfiability threshold can be computed analytically are indicated by *.

A	name	L_s	l_d	l_s
0100	1-in-3	3	2.256(3)	2.368(4)
01000	1-in-4	3	2.442(3)	2.657(4)
00100*	2-in-4	3	2.513	2.827
01010*	odd 4-PC	4	2.856	4
010000	1-in-5	3	2.594(3)	2.901(6)
001000	2-in-5	4	2.690(3)	3.180(6)
010100	1-or-3-in-5	5	3.068(3)	4.724(6)
010010	1-or-4-in-5	4	2.408(3)	3.155(6)

By computing the second moment $\langle Z^2 \rangle$ and using the Chebyshev's inequality, as in [24,25], we have shown that the annealed entropy is equal to the typical one; thus, the satisfiability threshold l_s is given by $s_{\text{ann}}(l_s) = 0$. Examples of LOPs for which this works are the parity checks $A = 01010$, as well as $A = 00100, 0001000, 0010100$, etc. Note that, for instance, $A = 010010$ does not belong to this class because its 0–1 symmetry is spontaneously broken.

Algorithms.—We attempt to find solutions to LOPs in their satisfiable phase using three algorithms which are among the best for hard random instances of the K -satisfiability problem: belief propagation decimation (Bpd) [14] (which is the same as survey propagation [13] in LOPs), stochastic local search (SLS) [26], and reinforced belief propagation (rBP) [27].

In Bpd, one uses the knowledge of marginal variable probabilities from BP equations in order to identify the most biased variable, fix it to its most probable value, and reduce the problem. In K -SAT, the SP decimation (which in LOPs is equivalent to Bpd) has been shown to be very efficient, on very large problems, even very near to the satisfiability threshold [13]. However, in LOPs, the BP decimation fails badly. For example, in the 1-or-3-in-5 SAT problem, on truncated Poisson graphs with $M = 2 \times 10^4$ constraints, the probability of success is about 25% at $\bar{l} = 2$, and less than 5% at already $\bar{l} = 2.3$, way below the clustering threshold $l_d \simeq 3.07$.

Although we do not know how to analyze directly the Bpd process, some mechanisms explaining the failure of the decimation strategy can be understood using the approach of [28]. The idea is to analyze a slightly simpler decimation process, where the variable to be fixed is chosen uniformly at random and its value is chosen according to its exact marginal probability, which is assumed to be approximated by BP. The reduced formula after θN steps is equivalent to the reduced formula created by choosing a solution uniformly at random and revealing a fraction θ of its variables. The number of variables which were either revealed or are directly implied by the revealed ones is denoted $\Phi(\theta)$. The performance of this “uniform” BP decimation can be understood from the shape of the function $\Phi(\theta)$, which we have computed from the cavity method.

In Fig. 1, we show that the theoretical curve $\Phi(\theta)$ agrees with numerical results in regular 1-or-3-in-5 SAT. At connectivity $L = 3$, the function has a discontinuity at $\theta_s \simeq 0.46$; thus, after fixing a fraction θ_s of variables, an infinite avalanche of direct implications follows and small errors in the BP estimation of marginals lead to a contradiction with high probability. At connectivity $L = 2$, the function $\Phi(\theta) \rightarrow 1$ at $\theta_1 \simeq 0.73$. This means that if a fraction $\theta > \theta_1$ of variables in a random solution is revealed, the residual problem has only this single solution. Any mistake in the previously fixed variables matters and causes a contradiction. In all the LOPs we have studied, $\Phi(\theta)$ has one these two fatal properties. The inset of Fig. 1 shows that, in

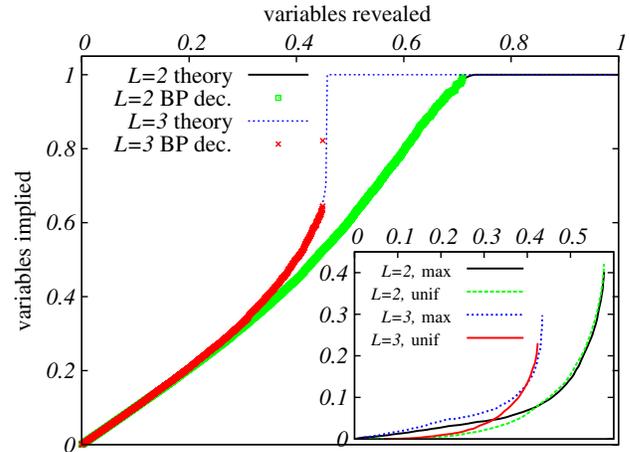


FIG. 1 (color online). Uniform BP decimation in regular 1-or-3-in-5 SAT with $L = 2$ and $L = 3$: plot of $\Phi(\theta)$, as obtained analytically (lines) and from the uniform BP decimation (points): the two plots agree perfectly. For $L = 3$, the decimation fails because of avalanches at the discontinuity of $\Phi(\theta)$; for $L = 2$, it fails when $\Phi(\theta) \rightarrow 1$ for $\theta < 1$. Inset: Comparison between Bpd and uniform BP decimation. The number of directly implied variables is plotted against number of variables which were free just before fixing them. The two methods are very close, and they fail at about the same value of θ .

LOPs, there is not much difference in the behaviors of Bpd and this uniform BP decimation.

Stochastic local search (SLS) algorithms exist in many different versions and are used in most practical cases where the exhaustive search is too time consuming. The main idea of the family of algorithms is to perform a random walk in configurational space, trying to minimize the number of unsatisfied constraints. In the implementation of [17], a variable which belongs to at least one unsatisfied constraint is chosen randomly. If flipping this variable does not increase the energy, the flip is accepted. If it increases the energy, the flip is accepted with probability p . This is repeated until either one finds a solution, or the number of steps per variable exceeds T . The parameter p must be optimized. In Fig. 2, we plot the fraction of successful runs for the 1-or-3-in-5 SAT with M constraints and $p = 0.00003$. Even with the largest value of T , we have not been able to solve instances with average connectivity larger than 3.05.

The belief propagation reinforcement (rBP) was originally introduced in [27]. The main idea is to add an external field $\mu_{s_i}^i$ which biases the variable i in the direction of the marginal probability computed from the BP messages. This modifies BP Eq. (2) to $\psi_{s_i}^{i \rightarrow a} = \mu_{s_i}^i \prod_{b \in \partial i - a} \psi_{s_b}^{b \rightarrow i} / Z^{i \rightarrow a}$. The algorithm then works as follows: Iterate the BP equations n -times. Update all the external fields: If $\xi_1^i < \xi_0^i$, set $\mu_1^i = \pi^{l_i}$, $\mu_0^i = (1 - \pi)^{l_i}$; otherwise, set $\mu_1^i = (1 - \pi)^{l_i}$, $\mu_0^i = \pi^{l_i}$ where $\xi_{s_i}^i = (\mu_{s_i}^i)^{1/(l_i - 1)} \prod_{a \in \partial i} \psi_1^{a \rightarrow i}$. At each iteration, one checks if the most probable configuration, given by $s_i = 0$ if $\mu_0^i > \mu_1^i$ and $s_i = 1$ otherwise, is a solution. If it is not, one

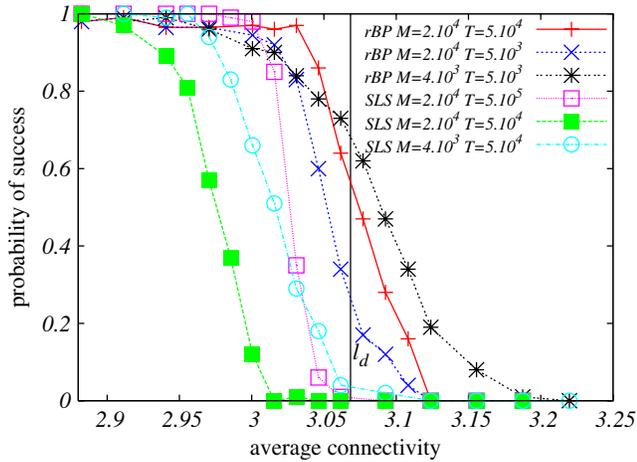


FIG. 2 (color online). Performance of reinforced BP and stochastic local search for 1-or-3-in-5 SAT with M constraints. The fraction of successful runs is plotted against the average connectivity \bar{l} . The clustering threshold l_d is marked, and the satisfiability transition is at $l_s = 4.72$. The maximal numbers of steps per variable T are chosen such that the running times of rBP and SLS are comparable.

iterates at maximum T times. We chose $n = 2$ and optimized the value of π . In Fig. 2, we plot the fraction of successful runs for the 1-or-3-in-5 SAT with $\pi = 0.42$ for $2.8 < \bar{l} < 3$ and $\pi = 0.43$ for $3 \leq \bar{l} < 3.2$. The performance is marginally better than SLS, but again one cannot penetrate into the clustered phase.

We have observed the same behavior for all LOPs we studied: the clustering transition point l_d seems to be a boundary beyond which all these three algorithms fail. As shown in Table I, this point can be very far from the SAT-UNSAT transition l_s , meaning that there is a broad range of instances where known algorithms are totally inefficient. The parity check problems are the exception as they can be solved with linear programming algorithms.

Conclusions—LOPs make a broad class of extremely hard constraint satisfaction problems. Their phase diagram is simple: the set of satisfiable configurations becomes clustered when the average connectivity is $\bar{l} > l_d$, and it disappears for $\bar{l} > l_s$. These two thresholds can be computed efficiently using population dynamics, and in the case of some symmetric problems, the value of l_s can be confirmed rigorously. At the same time, the best algorithms known for random CSP fail to find solutions in the clustered phase $l_d < \bar{l} < l_s$. This difficulty is due to the “locked” nature of the problem which reduces the clusters to single points. It will be interesting to investigate if LOPs might be used to design new efficient nonlinear error-correcting codes, or if the planted LOPs are good candidates for one-way functions in cryptography.

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